# Inorganic Chemistry

## $F^+$ and $F^-$ Affinities of Simple N<sub>x</sub>F<sub>y</sub> and O<sub>x</sub>F<sub>y</sub> Compounds

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Atomization energies at 0 K and heats of formation at 0 and 298 K are predicted for the neutral and ionic N<sub>x</sub>F<sub>v</sub> and O<sub>x</sub>F<sub>v</sub> systems using coupled cluster theory with single and double excitations and including a perturbative triples correction (CCSD(T)) method with correlation consistent basis sets extrapolated to the complete basis set (CBS) limit. To achieve near chemical accuracy ( $\pm 1$  kcal/mol), three corrections to the electronic energy were added to the frozen core CCSD(T)/CBS binding energies: corrections for core-valence, scalar relativistic, and first order atomic spin-orbit effects. Vibrational zero point energies were computed at the CCSD(T) level of theory where possible. The calculated heats of formation are in good agreement with the available experimental values, except for FOOF because of the neglect of higher order correlation corrections. The F<sup>+</sup> affinity in the N<sub>x</sub>F<sub>y</sub> series increases from N<sub>2</sub> to N<sub>2</sub>F<sub>4</sub> by 63 kcal/mol, while that in the O<sub>2</sub>F<sub>y</sub> series decreases by 18 kcal/mol from O<sub>2</sub> to O<sub>2</sub>F<sub>2</sub>. Neither N<sub>2</sub> nor N<sub>2</sub>F<sub>4</sub> is predicted to bind  $F^-$ , and  $N_2F_2$  is a very weak Lewis acid with an  $F^-$  affinity of about 10 kcal/mol for either the cis or trans isomer. The low F<sup>-</sup> affinities of the nitrogen fluorides explain why, in spite of the fact that many stable nitrogen fluoride cations are known, no nitrogen fluoride anions have been isolated so far. For example, the F<sup>-</sup> affinity of NF is predicted to be only 12.5 kcal/mol which explains the numerous experimental failures to prepare NF2<sup>-</sup> salts from the well-known strong acid HNF<sub>2</sub>. The F<sup>-</sup> affinity of O<sub>2</sub> is predicted to have a small positive value and increases for O<sub>2</sub>F<sub>2</sub> by 23 kcal/ mol, indicating that the O<sub>2</sub>F<sub>3</sub><sup>-</sup> anion might be marginally stable at subambient temperatures. The calculated adiabatic ionization potentials and electron affinities are in good agreement with experiment considering that many of the experimental values are for vertical processes.

#### Introduction

Compounds with fluorine atoms adjacent to other atoms with lone pairs pose problems for computational chemistry. However, one can make reliable estimates of many of their thermomdynamic properties by the judicious use of various approaches. We are especially interested in the  $F^+$  and  $F^-$  affinities<sup>1-6</sup> of a broad range of compounds to make reac-

tivity predictions. Given the limited experimental data available for the fluorides of nitrogen and oxygen, we have performed high accuracy theoretical calculations to determine thermodynamic properties of these molecules, namely, their heats of formation and  $F^+$  and  $F^-$  affinities.

There have been numerous experimental investigations of nitrogen-fluorine and oxygen-fluorine compounds, and Klapötke recently reviewed several classes of nitrogen-fluorine compounds.<sup>7</sup> The experimental heats of formation of NF,<sup>8</sup> NF<sub>2</sub>,<sup>8</sup> NF<sub>3</sub>,<sup>8</sup> *cis*- and *trans*-N<sub>2</sub>F<sub>2</sub>,<sup>8</sup> N<sub>2</sub>F<sub>4</sub>,<sup>8</sup> and O<sub>2</sub>F<sub>2</sub>,<sup>9,10</sup> are available. The vibrational frequencies of NF<sub>2</sub>,<sup>11,12</sup> HNF<sub>2</sub>

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(included for its acidity),  ${}^{13}NF_{3,7}{}^{13}NF_{4}{}^{+14}$  trans-FN= NF,  ${}^{13,15}$  cis-FN=NF,  ${}^{16}N_2F^+$ ,  ${}^{17}N_2F_3^+$ ,  ${}^{18}N_2F_4$ ,  ${}^{19}$  and FN<sub>3</sub><sup>20</sup> have been reported. The molecular structures of NF<sub>2</sub>,  ${}^{21}$  cis-<sup>22,23</sup> and trans-N<sub>2</sub>F<sub>2</sub>,  ${}^{22}$  and gauche<sup>21,24,25</sup> and trans-N<sub>2</sub>F<sub>4</sub><sup>21,24,25</sup> based on electron diffraction studies have been reported. For gaseous<sup>26</sup> and liquid<sup>27,28</sup> N<sub>2</sub>F<sub>4</sub>, infrared spectroscopic experiments have indicated the existence of an equilibrium mixture of two rotational isomers, trans ( $C_{2h}$ symmetry) and gauche ( $C_2$  symmetry), that differ slightly in energy. The conformational stability of gaseous N<sub>2</sub>F<sub>4</sub> has been studied by far-infrared and low-frequency Raman spectroscopy.<sup>19</sup> The molecular structure of  $N_2F_4$  has been reported from microwave rotational spectroscopy studies.<sup>29,30</sup> The rotational spectrum of the  $NF_2$  free radical in the millimeter-wave region<sup>31</sup> and its microwave<sup>32</sup> and absorption<sup>33</sup> spectra have been reported. The photodissociation of  $NF_3$  has been studied by vacuum-ultraviolet fluorescence spectroscopy.<sup>34</sup> The crystal structures of several  $N_2F^+$  and  $NF_4^+$  salts have been reported.<sup>17,35,36</sup> The structure of the free gaseous ion  $N_2F^+$  has also been determined using millimeter-wave spectroscopy.<sup>37</sup> The experimental structure<sup>38,39</sup>

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and anharmonic vibrational frequencies<sup>40-42</sup> of FOOF have also been reported.

There have been several theoretical investigations into some of the compounds under study. FOOF has proven to be a troublesome molecule for traditional single reference ab initio methods, because of the presence of a large number of nominally inactive lone pairs that can interact with each other at short distances. Thus, several theoretical approaches have been employed at describing FOOF accurately.<sup>43-45</sup> The heats of formation and structures of five small oxygen fluoride molecules have been determined with coupled cluster theory incorporating at least quadruple excitations given that the wave functions of several systems have large multiconfiguration character requiring correlation recovery beyond CCSD(T) to achieve accurate results.<sup>43</sup> The  $\Delta H_{f,298K}$  of FOOF was predicted to be  $6.4 \pm 0.7$  kcal/mol, which is the best available computational value so far.43 Feller et al. previously reported coupled cluster and multireference configuration interaction calculations of FOOF.44,45 Martin and co-workers have also reported W4.n calculations on the fluorine oxides.<sup>46</sup> They performed CCSDT(Q)/cc-pVDZ calculations to estimate the higher order correlation corrections for the geometry of FOOF. Their best value for  $\Delta H_{f,298\rm K}(\rm FOOF)$  at the W4 level is 7.84  $\pm$  0.18 kcal/mol using the CCSD(T)/cc-pVQZ geometry. FOOF was also studied using the analytic gradient technique at the CCSD-(T)(FC)/[5s,3p,3d,1f] level.<sup>47</sup> Lee et al.<sup>48</sup> reported  $\Delta H_{f,0K}$ - $(FOOF) = 9.7 \pm 2.0$  kcal/mol at the CCSD(T)(FC) level with a triple- $\zeta$  atomic natural orbital basis set. Ventura and Keininger studied FOOF with B3LYP/6-311++G(3df,3pd)density functional theory (DFT) obtaining  $\Delta H_{f,0K}(FOOF) =$ 8.2 kcal/mol based on an average of two procedures.<sup>49</sup> The geometry of FOOF was reported at the DFT level within the local density approximation (LDA).<sup>50</sup> Several DFT methods were used to investigate the structures FOOF and three other small oxygen fluorides.<sup>51</sup> Keininger et al.<sup>52</sup> studied a collection of oxygen fluoride molecules at the DFT-B3PW91 level with the cc-pVQZ or aug-cc-pVQZ basis sets, and report  $\Delta H_{f,298K}(FOOF) = 7.3$  and 4.1 kcal/mol at the DFT and Gaussian-2 levels, respectively. Kraka et al. reported the geometry of FOOF at the CCSD(T)/CBS level based on extrapolating the values obtained with the aug-cc-pVTZ and aug-cc-pVQZ basis sets.5

Zhang and co-workers have studied the low-lying electronic states of NF2 at the MP2 level.<sup>54</sup> High level calculations at the CCSD(T)/CBS level on the structures and thermodynamic properties of  $N_2F^+$ ,  $N_2F_2$  (cis and trans), and  $F_2NN$ have recently been reported,<sup>36</sup> and the results are consistent

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with other high quality calculations.<sup>37,55</sup> The electron affinity of NF has been calculated at R(U)CCSD(T), icMRCI(+Q), and MR-ACPF levels with the (d)-aug-cc-pVnZ basis sets up to n = 6.56 The calculated value of 0.247 eV is 0.25 eV lower than the reported experimental value of 0.5 eV; however, no experimental details were given.<sup>57</sup>

#### **Computational Methods**

Modern computational chemistry methods can now provide reliable predictions of thermodynamic properties to within about 1 kcal/mol for most compounds that are not dominated by multireference character.58 We use the approach that we have been developing at The University of Alabama and Washington State University for the prediction of accurate molecular thermochemistry to determine the atomization energies and the heats of formation of these compounds.<sup>43–45</sup> Our approach is based on calculating the total atomization energy (TAE) of a molecule and using this value with known heats of formation of the atoms to calculate the heat of formation at 0 K. The approach starts with coupled cluster theory with single and double excitations and including a perturbative triples correction (CCSD-(T)), $^{59-61}$  combined with the correlation-consistent basis sets<sup>62</sup> extrapolated to the complete basis set (CBS) limit to treat the correlation energy of the valence electrons. This is followed by a number of smaller additive corrections including core-valence and relativistic effects, both scalar and spin-orbit. The zero point energy can be obtained from experiment, theory, or a combination of the two. Corrections to 298 K can then be calculated by using standard thermodynamic and statistical mechanics expressions in the rigid rotor-harmonic oscillator approximation<sup>63</sup> and appropriate corrections for the heat of formation of the atoms.

The standard aug-cc-pVnZ basis sets were used for N, O, and F and abbreviated as aVnZ. Only the spherical component subsets (e.g., 5-term d functions, 7-term f functions, etc.) of the Cartesian polarization functions were used. All CCSD-(T) calculations were performed with the MOLPRO-2006<sup>65</sup> program system on an SGI Altix, a Cray XD-1, or the dense memory Linux cluster at the Alabama Supercomputer Center, or on the Dell Linux cluster at The University of Alabama, or on the massively parallel HP Linux cluster in the Molecular Science Computing Facility (MSCF) in the

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For the open shell atomic calculations, we used the restricted method for the starting Hartree-Fock wave function and then relaxed the spin restriction in the coupled cluster portion of the calculation. This method is conventionally labeled R/UCCSD(T).<sup>66–68</sup> Our CBS estimates use a mixed exponential/Gaussian function of the form<sup>69</sup>

$$E(n) = E_{\text{CBS}} + Be^{-(n-1)} + Ce^{-(n-1)^2}$$
(1)

where n = 2 (aVDZ), 3 (aVTZ), and 4 (aVQZ).

To achieve thermochemical properties within  $\pm 1$  kcal/mol of experiment, it is necessary to account for core-valence correlation energy effects beyond whatever treatment is used for the valence electron correlation. Core-valence (CV) calculations were carried out with the weighted CV basis set cc-pwCVTZ.70 Relativistic effects are included at the scalar relativistic level,  $\Delta E_{\rm SR}$ , plus the spin orbit corrections for atoms. The atomic spin-orbit corrections are  $\Delta E_{SO}(O) =$ 0.22,  $\Delta E_{SO}(F) = 0.39$ , and  $\Delta E_{SO}(F^+) = 0.48$  kcal/mol, respectively, taken from the tables of Moore.<sup>71</sup> We evaluated  $\Delta E_{\rm SR}$  by using expectation values for the two dominant terms in the Breit-Pauli Hamiltonian, the so-called mass-velocity and one-electron Darwin (MVD) corrections from configuration interaction singles and doubles (CISD) calculations. The quantity  $\Delta E_{SR}$  was obtained from the CISD wave function with aVTZ basis set at the appropriate optimized geometry.

The geometries of the diatomic molecules were optimized up through the CCSD(T)/aVQZ level. Bond distances, harmonic frequencies, and anharmonic constants for the diatomics were obtained from a fifth order Dunham fit72 of the potential energy curve at the CCSD(T)/aVQZ level. For  $NF_2^{0/\pm}$ ,  $HNF_2$ ,  $NF_3^{0/\pm}$ ,  $NF_4^{\pm}$ ,  $NH_5$ ,  $NF_5$ ,  $N_2F^{0/\pm}$ ,  $N_2F_2$ isomers,  $FN_3$ ,  $F_2N_3^{\pm}$  (except for  $F_2N_3^{+}$  (B,  $C_1$ ), which was optimized at the MP2/aVTZ level),  $OF_2$ , and  $OF_3^{\pm}$ , the geometries were optimized up through the CCSD(T)/aVTZ level. The aVTZ geometries were then used in single point CCSD(T)/aVQZ calculations. Zero point energies ( $\Delta E_{ZPE}$ ) were calculated at the CCSD(T)/aVTZ level without scaling. For FN<sub>3</sub>,  $\Delta E_{ZPE}$  was calculated at the CCSD(T)/VTZ level without scaling, and for the  $F_2N_3^{\pm}$  ions at the MP2/aVTZ level. For the remaining molecules, geometry optimizations were performed and  $\Delta E_{\rm ZPE}$ 's were calculated at the MP2/ aVTZ level; this geometry was used in single point CCSD(T) calculations with the aVDZ, aVTZ, and aVQZ basis sets for these molecules.

By combining our computed  $\sum D_0$  values, given by the following expression,

$$\sum D_0 = \Delta E_{\text{elec}}(\text{CBS}) - \Delta E_{\text{ZPE}} + \Delta E_{\text{CV}} + \Delta E_{\text{SR}} + \Delta E_{\text{SO}}$$
(2)

with the known heats of formation at 0 K for the elements, we can derive  $\Delta H_{f,0K}$  values for the molecules under study. The heats of formation of H, N, O, and F are well-established as

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**Table 1.** Optimized MP2/aVTZ and CCSD(T) Bond Lengths (Å) and Bond Angles  $(deg)^{a,b}$ 

molecule	$R_{ m NF/OF}$	$R_{ m NN/OO/FF}$	∠FNN/FOO	∠FNF/FOF/NNN	dihedral
$NF^{+}(^{2}\Pi, C_{\infty\nu})$ $NF(^{3}\Sigma^{-}, C_{\infty\nu})$	1.1854				
$N\Gamma(2, C_{\infty \nu})$	1.317 <sup>76</sup>				
$NF^{-}(^{2}\Pi, C_{\infty v})$	1.4919				
$OF^{+}(\Sigma, C_{\infty y})$ $OF^{2}\Pi(C_{-})$	1.2337 1.3536 <sup>43</sup>				
$OI (II, C_{\infty \nu})$	1.3541 <sup>84</sup>				
$OF^{-}(\Sigma^{+}, C_{\infty v})$	1.5170				
$F_2^+({}^{2}\Pi_g, D_{\infty h})$		1.3075			
$\Gamma_2(\Sigma_g, D_{\infty h})$		1.4130 1 4119 <sup>76</sup>			
$F_2^{-}(^2\Sigma_u^{+}, D_{\infty h})$		1.9225			
$NF_2^+({}^1A_1, C_{2\nu})$	1.2504			107.8	
$NF_2(^2B_1, C_{2\nu})$	1.3536			103.1 $103.1 \pm 0.02^{31}$	
	$1.363 \pm 0.008^{21}$			$102.5 \pm 0.9^{21}$	
$NF_2^{-}(^{1}A_1, C_{2\nu})$	1.4965			99.4	
$HNF_2({}^{1}A', C_s)$	1.0272 (1,2)			103.0	102.5 (2,1,3,4)
$NF_2^+ (^2A_1, C_2)$	1.3997 (1,3) 1.2858			113 7	132 3 (2 1 4 3)
$NF_3({}^{1}A_1, C_{3v})$	1.3702			101.8	10210 (2,1,1,0)
	1.371 <sup>74,75</sup>			102.9 <sup>74,75</sup>	
$NF_3$ ('A', $C_s$ )	1.3823(1,3) 2 2391(1,2)			101.6(3,1,4) 96.9(2,1,3)	98.5 (2,1,3,4)
$NF_4^+$ ( <sup>1</sup> A <sub>a</sub> , $T_d$ )	1.3112			109.5	120.0
	$1.3076^{35}$			$109.5^{35}$	
$NF_4(^2A_1, C_{3v})$	2.7017 (1,3)			101.9(2,1,4)	120.0 (2,1,3,4)
$NE_4^{-1}(^1A_1, C_{2n})$	1.3688(1,2) 1.800(1,2)			116.2(2,1,3) 171.0(2.1.4)	-1731(2134)
(11, 020)	1.3720 (1,3)			101.8 (3,1,5)	-93.5 (2,3,1,5)
MIT (LA / D )	1 (770)			92.8 (2,1,3)	120.0
$NH_5({}^{t}A_1{}', D_{3h})$	$1.6778_{ax}$ 1.0006			90.0	120.0
$NF_5({}^{1}A_1', D_{3h})$	1.5778 <sub>ax</sub>			90.0	120.0
1	1.3817 <sub>eq</sub>				
$N_2({}^{1}\Sigma^{+}_{g}, D_{\infty h})$		1.1005			
$N_2 F^+ ({}^1\Sigma_2 C_{-1})$	1 2357	1.0977	180.0		
$1\sqrt{2}$ ( $-g$ , $c_{\infty p}$ )	1.217 <sup>17</sup>	1.099 <sup>17</sup>	10010		
$N_2 F(^2 \Sigma_g, C_{\infty \nu})$	2.8263	1.1138	180.0		
$N_2F$ ('A', $C_s$ )	2.8759	1.1143	/9.6		
$cis-N_2F_2$ ( <sup>1</sup> A <sub>1</sub> , $C_{2\nu}$ )	1.3872	1.2248	114.2		0.0
	$1.410 \pm 0.009^{22}$	$1.214 \pm 0.008^{22}$	$114.4 \pm 1.0^{22}$		
trans N E $(^{1}\Lambda C)$	$1.384 \pm 0.010^{23}$	$1.214 \pm 0.005^{23}$	$114.5 \pm 0.5^{23}$		180.0
$II ull s - IN_2 I'_2 (A_g, C_{2h})$	1.3860 $1.396 \pm 0.008^{22}$	1.2321 $1.231 \pm 0.010^{22}$	104.7 $105.5 \pm 0.7^{22}$		180.0
$N_2F_3^+({}^1A', C_s)$	1.2880 (1,4)	1.2526	127.4 (4,1,2)	113.6 (4,1,5)	0.0 (3,2,1,4)
	1.2836 (1,5)		119.0 (5,1,2)		180.0 (3,2,1,5)
	1.3140 (2,3)		108.2(3.2.1)		
$N_2F_3(^2A, C_1)$	1.3853 (1,4)	1.4188	100.2(3,2,1) 102.2(3,2,1)	103.1 (4,1,5)	176.1 (3,2,1,5)
2 3 ( ) 3	1.3755 (1,5)		105.2 (4,1,2)		69.2 (3,2,1,4)
$\mathbf{N} \mathbf{E} = (1 \mathbf{A} \cdot \mathbf{C})$	1.3440 (2,3)	1 2412	100.6 (5,1,2)	010(415)	00.4(2.2.1.4)
$N_2F_3$ (A, C <sub>1</sub> )	2.2192 (1,4) 1 3943 (1.5)	1.2413	103.8(3,2,1) 102.1(4.1,2)	91.9 (4,1,5)	-90.4(3,2,1,4) 174.7(3.2,1,5)
	1.4073 (2,3)		103.3 (5,1,2)		1, 11, (0,2,1,0)
$N_2F_4({}^1A_g, C_{2h})$	1.3741	1.5016	99.8	102.8	75.0 (3,2,1,6)
$N_{\rm e} F_{\rm e} \begin{pmatrix} 1 \Delta & C_{\rm e} \end{pmatrix}$	$1.375 \pm 0.004^{24}$ 1.3767 (1.5)	$1.489 \pm 0.007^{24}$ 1.4628	$100.6 \pm 0.6^{24}$ 100.7 (3.2.1)	$102.9 \pm 0.75^{24}$ 102.9 (3.2.4)	-1741(3216)
(1, 2, 2, 3, 3, 3, 3, 3, 3, 3, 3, 3, 3, 3, 3, 3,	1.3737 (1,6)	1.4020	107.1 (4,2,1)	102.9 (3,2,4)	40.3 (4,2,1,5)
	$1.375 \pm 0.004^{24}$	$1.489 \pm 0.007^{24}$	$100.1 \pm 1^{24}$	$105.1 \pm 1^{24}$	( ) / / / /
$\mathbf{N} \mathbf{E}^{\pm} (\mathbf{A} \mathbf{A} \mathbf{C})$	$1.393 \pm 0.008^{21}$	$1.53 \pm 0.02^{21}$	$104.3 \pm 1^{21}$	$103.7 \pm 0.9^{21}$	1744(41(2))
$1N_2\Gamma_5$ (A, $C_s$ )	1.3130(0,7) 1.3084(6.2)	1.0304	107.8 (3,0,1) 99.9 (4 1 6)	109.2 (2,0,7)	-67.4(4,1,6,2)
	1.3280 (1,4)		133.3 (7,6,1)	104.7 (4,1,5)	53.5 (4,1,6,7)
$N_2F_5^{-}(^1A, C_1)$	1.4876 (1,4)	3.0258	64.3 (4,1,6)	99.4 (4,1,5)	-82.1 (4,1,6,3)
	1.4796 (1,5)		80.4 (5,1,6)	100.5 (3,6,2)	1/5.8(4,1,6,2) 187(4167)
	1.3589 (6.3)		86.0 (3.6.1)	101.5 (3.6.7)	172.5 (5.1.6.3)
	1.3773 (6,7)		115.0 (7,6,1)		70.4 (5,1,6,2)
DI du a	1.400.4	1 1075 (1 2)	102.0	170 1	-86.8 (5,1,6,7)
$r_{1N_3}(A, C_s)$	1.4294	1.13/3(1,2)	103.9	1/2.1	180.0

Table 1. Continued	Table	1.	Continue	d
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molecule	$R_{ m NF/OF}$	$R_{ m NN/OO/FF}$	∠FNN/FOO	$\angle$ FNF/FOF/NNN	dihedral
	z ( , 20	1.2671 (1,3)	100.020	1=0.020	
	1.44420	$1.132(1,2)^{20}$ 1.252(1,2) <sup>20</sup>	103.820	170.920	
$\mathbf{F} \mathbf{N} + (1\mathbf{A}' \mathbf{C} \mathbf{A})$	1 2550	$1.233(1,3)^{-3}$ 1.1026 (1.2)	07.6	107.4	125.6(4.2,1.2)
$\Gamma_2 \Gamma_3$ (A, C <sub>s</sub> , A)	1.2350	2,4677(1,2)	97.0	107.4	123.0 (4,2,1,3)
$F_2N_2^+$ ( <sup>1</sup> A C <sub>1</sub> B)	1 2977 (2 5)	1.2051(1.2)	1353(521)	87 3 (1 2 3)	179 3 (5 2 3 4)
1 21 (3 (11, 01, D)	1 3254 (3 4)	1 4221 (2.3)	105.2(4.3.2)	07.5 (1,2,5)	-98.6(4.3.2.1)
$F_2N_3^+$ ( <sup>1</sup> A, C <sub>2</sub> , C)	1.3293	1.2251	111.2	152.7	135.8
$F_2N_3^{-}({}^1A', C_s, B)$	2.0019 (1.3)	1.2851 (1.2)	98.8 (3.1.2)	160.6	180.0 (5.2.1.3)
2 5 ( ) - 37 )	1.5298 (2,5)	1.1409 (2,4)	86.0 (5,2,1)		0.0(3,1,2,4)
$F_2N_3^{-}(^1A', C_s, C)$	1.5082 (2,4)	1.2708 (1,2)	103.0 (4,2,1)	129.9	180.0 (4,2,1,3)
	1.6082 (3,5)	1.2330 (1,3)	108.6 (5,3,1)		0.0 (5,3,1,2)
$OF_2(^1A_1, C_{2\nu})$	1.4033			102.9	
	$1.412^{85}$			103.1 <sup>85</sup>	
$OF_3^+(^1A_1, C_{3\nu})$	1.3768			103.5	107.8 (4,1,3,2)
$OF_3^{-}({}^{1}A_1, C_{2\nu})$	1.4567 (1,3)			100.3 (2,1,3)	180.0 (2,1,3,4)
	1.7432 (1,2)			159.5 (2,1,4)	
$O_2F^+({}^1A', C_s)$	1.4639	1.1291	112.9		
$O_2F^+({}^{3}A'', C_s)$	1.3270	1.2589	108.6		
$O_2F(^2A^{\prime\prime},C_s)$	1.641443	1.195943	110.943		
2	$1.649^{83}$	$1.200^{83}$	$111.2^{83}$		
$O_2F^-({}^{3}A'', C_s)$	2.8437	1.2275	177.6		
$O_2F^-(A', C_s)$	1.9662	1.2655	109.4		
$O_2F_2(^{1}A, C_2)$	1.6188	1.1662	110.5	104.4	-88.7 (3,2,1,4)
	$1.575 \pm 0.003^{38}$	$1.217 \pm 0.003^{38}$		$109.5 \pm 0.5^{38}$	
	$1.586 \pm 0.002^{39}$	$1.216 \pm 0.002^{39}$		$109.2 \pm 0.2^{39}$	
$O_2F_3^+({}^{1}A', C_s)$	1.3605 (1,4)	1.6533	104.6 (4,1,2)	103.2 (3,2,5)	-54.6 (3,2,1,4)
	1.4054 (2,5)		103.7 (3,2,1)		
$O_2F_3^{-}({}^{^{t}}A', C_s)$	1.8386 (2,3)	1.1866	111.8 (3,2,1)	86.0 (3,2,5)	-132.7 (3,2,1,4)
	2.0733(1.4)		129.0(4.1.2)		

<sup>a</sup> Experimental values are given in italics. <sup>b</sup> Geometrical parameters for NF<sub>2</sub><sup>0/±</sup>, HNF<sub>2</sub>, NF<sub>3</sub><sup>0/±</sup>, NF<sub>4</sub><sup>±</sup>, NH<sub>5</sub>, NF<sub>5</sub>, N<sub>2</sub>F<sup>0/±</sup>, N<sub>2</sub>F<sub>2</sub> isomers, FN<sub>3</sub>, F<sub>2</sub>N<sub>3</sub><sup>±</sup>,  $OF_2$ , and  $OF_3^{\pm}$  were calculated at the CCSD(T)/aVTZ level and for the diatomics at the CCSD(T)/aVQZ level.

 $\Delta H_{f,0K}(H) = 51.63 \text{ kcal/mol}, \Delta H_{f,0K}(B) = 135.1 \text{ kcal/mol},^{73}$  $\Delta H_{f,0K}(N) = 112.53 \text{ kcal/mol}, \\ \Delta H_{f,0K}(O) = 58.99 \text{ kcal/mol},$ and  $\Delta H_{f,0K}(F) = 18.47 \text{ kcal/mol.}^8$  Heats of formation at 298 K were obtained by following the procedures outlined by Curtiss et al.<sup>64</sup>

#### **Results and Discussion**

The calculated and experimental geometries for the molecules under study are given in Table 1, including their electronic states and symmetry labels. The optimized molecular structures are depicted in Figure 1. The total CCSD(T) energies and calculated CCSD(T) or MP2 harmonic frequencies for the molecules are given as Supporting Information in Table SI-1 and Table SI-2, respectively, where they are compared with the available experimental values. The component reaction energies are also provided in the Supporting Information (Table SI-3).

Geometries. The experimental geometry of  $NF_2$  has been reported from rotational spectroscopy<sup>31</sup> and electron diffraction experiments,<sup>21</sup> and our CCSD(T)/aVTZ value for the r(NF) distance is in excellent agreement, within |0.009| Å, of the two experimental values. Our CCSD(T)/aVTZ geometry of NF<sub>3</sub> is in good agreement with the experimental structures from microwave<sup>74</sup> and electron diffraction studies.<sup>75</sup>

The cationic and anionic forms of N<sub>2</sub>F exhibit very different structures with  $N_2F^+$  predicted to be linear<sup>3</sup> and N<sub>2</sub>F<sup>-</sup> predicted to be bent with an  $\angle$  FNN of  $\sim 80^{\circ}$ .

The r(NN) distance of N<sub>2</sub>F<sup>-</sup> is only slightly longer than that in  $N_2 ({}^{1}\Sigma_g^{+})^{76}$  and slightly shorter than that in  $N_2F^+$ . The the r(NF) distance is longer by 1.640 Å in N<sub>2</sub>F<sup>-</sup> than in  $N_2F^+$  so,  $N_2F^-$  is best considered as a weak complex of  $F^-$  with  $N_2$ .

The structures of  $N_2F_3^+$  and  $N_2F_3^-$  also differ significantly with  $N_2F_3^{\,+}$  being planar with 3 tightly bound fluorines and  $N_2F_3^{\,-}$  being nonplanar with one very loosely bound fluorine. As predicted for the  $N_2F^{\pm}$  pair, the r(NN) distance in the anion is predicted to be slightly shorter than in the cation. The values for r(NN) in N<sub>2</sub>F<sub>3</sub><sup>+</sup> and  $N_2F_3^-$  are typical N=N double bond values and are only slightly longer than that in *trans*- $N_2F_2$ . Similarly, the r(NF) distances in the N<sub>2</sub>F<sub>2</sub> subunit are predicted to be only slightly longer than that in *trans*-N  $_2F_2$ , but the  $N_2F_2$ --F<sup>-</sup> distance is again very long (2.219 Å). Therefore, the structure of  $N_2F_3^-$  can also be considered as a weak complex of  $F^-$  loosely interacting with *trans*-N<sub>2</sub>F<sub>2</sub> from above the plane of the molecule.

There are two rotational isomers of  $N_2F_4$ , gauche ( $C_2$ symmetry) and trans ( $C_{2h}$  symmetry), that are close in energy.<sup>24,27,28</sup> Bauer and Cardillo<sup>24</sup> determined the trans rotamer to be more stable than the gauche rotamer by 0.3–0.5 kcal/mol with a 47% gauche–53% trans mixture below room temperature. At the CCSD(T)/CBS level plus the additional corrections, we predict the trans rotamer to be more stable than the gauche rotamer by 0.1 kcal/mol at 298 K, completely consistent with the experimental

<sup>(73)</sup> Karton, A.; Martin, J. M. L. J. Phys. Chem. A 2007, 111, 5936.
(74) Sheridan, J.; Gordy, W. Phys. Rev. 1950, 79, 513.
(75) Schomaker, V.; Lu, C. J. Am. Chem. Soc. 1950, 72, 1182.

<sup>(76)</sup> Huber, K. P.; Herzberg, G. Constants of Diatomic Molecules. Molecular Spectra and Molecular Structure, Vol. IV, Van Nostrand: Princeton, 1979





**Figure 1.** Optimized molecular structures for NF<sub>2</sub><sup>+</sup>, NF<sub>2</sub>, NF<sub>2</sub><sup>-</sup>, NF<sub>3</sub><sup>+</sup>, NF<sub>3</sub>, NF<sub>3</sub><sup>-</sup>, NF<sub>4</sub><sup>+</sup>, NF<sub>4</sub>, NF<sub>4</sub><sup>-</sup>, N<sub>2</sub>F<sup>+</sup>, N<sub>2</sub>F, N<sub>2</sub>F<sup>-</sup>, *cis*-N<sub>2</sub>F<sub>2</sub>, *trans*-N<sub>2</sub>F<sub>2</sub>, N<sub>2</sub>F<sub>3</sub><sup>+</sup>, N<sub>2</sub>F<sub>3</sub>, N<sub>2</sub>F<sub>3</sub><sup>-</sup>, N<sub>2</sub>F<sub>4</sub>, N<sub>2</sub>F<sub>5</sub><sup>-</sup>, N<sub>3</sub>F, N<sub>3</sub>F<sub>2</sub><sup>+</sup>, N<sub>3</sub>F<sub>2</sub><sup>-</sup>, NF<sub>5</sub>, OF<sub>2</sub>, OF<sub>3</sub><sup>+</sup>, OF<sub>3</sub><sup>-</sup>, O<sub>2</sub>F<sup>+</sup>, O<sub>2</sub>F<sup>-</sup>, O<sub>2</sub>F<sub>2</sub>, O<sub>2</sub>F<sub>3</sub><sup>+</sup>, and O<sub>2</sub>F<sub>3</sub><sup>-</sup>.

product ratio. At just the CCSD(T)/CBS level based on the valence electronic energies, the trans rotamer is more stable than the gauche by 0.2 kcal/mol. The experimental structure of the gauche<sup>21,24</sup> and trans<sup>24</sup> rotamers have been reported from electron diffraction studies. The MP2/aVTZ geometry parameters for the trans rotamer are in excellent agreement with the most recent electron diffraction structure,<sup>24</sup> with the largest discrepancy predicted for the *r*(NN) distance, which was calculated to be too long by 0.013 Å. Reasonable agreement is also found for the gauche rotamer with the r(NN) distance predicted to be shorter by 0.026 Å. The large difference in r(NN)between N<sub>2</sub>F<sub>4</sub> ( $C_{2h}$ ) (1.502 Å) and *trans*-N<sub>2</sub>F<sub>2</sub> (1.232 Å) is due to N<sub>2</sub>F<sub>4</sub> having an N–N single bond and N<sub>2</sub>F<sub>2</sub> having an N=N double bond. The r(NF) distances in N<sub>2</sub>F<sub>4</sub> ( $C_{2h}$ ) and N<sub>2</sub>F<sub>2</sub> are within 0.006 Å of each other.

Addition of F<sup>+</sup> and F<sup>-</sup> to N<sub>2</sub>F<sub>4</sub> gives the N<sub>2</sub>F<sub>5</sub><sup>+</sup> and N<sub>2</sub>F<sub>5</sub><sup>-</sup> ions with  $C_s$  and  $C_1$  symmetry structures, respectively. Whereas the N-F bonds of ~1.31 Å in N<sub>2</sub>F<sub>5</sub><sup>+</sup> are relatively short and stable, the r(NN) distance of 1.650 Å

is relatively long, implying that  $N_2F_5^+$  might at best have marginal stability at low temperatures. For example,  $N_2F_4$  with an N–N bond length of 1.502 Å dissociates already at room temperature to NF<sub>2</sub> radicals. The predicted instability of the  $N_2F_5^+$  cation is in accord with previous literature reports. In 1976, Toy and Stringham<sup>77</sup> reported the isolation of a stable solid which they attributed to  $[N_2F_5]^+[(CF_3)_3CO]^-$  but which was subsequently reassigned by Christe and co-workers<sup>78</sup> to  $(NO)_2SiF_6$ . In a second report,<sup>79</sup> Stringham and Toy also claimed the synthesis of  $N_2F_5^+BF_4^-$  which was again refuted by Christe and co-workers. By analogy with  $N_2F^-$  and  $N_2F_3^-$ , the  $N_2F_5^-$  anion is also unstable. Its predicted N-N bond distance of 3.026 Å is only slightly shorter than the sum of its van der Waals radii  $(r_{VDW}(N) = 1.55 \text{ Å}).^{80}$ 

There has been substantial interest in the existence of hypercoordinated nitrogen(V) compounds with more than three substituents. We predict NF<sub>5</sub>  $(D_{3h})$  to be a local minimum from a harmonic vibrational analysis at the CCSD(T)/aVTZ level, in agreement with the results of Bettinger et al.<sup>81</sup> who found a minimum at all levels of theory employed, up through CCSD(T)/DZP. However, as discussed below in more detail, NF<sub>5</sub> is thermodynamically unstable by 42 kcal/mol with respect to decomposition to NF<sub>3</sub> and F<sub>2</sub>, and the barrier to the loss of an F atom is only about 16 kcal/mol.<sup>72</sup>

The structure of FN<sub>3</sub> is planar with  $C_s$  symmetry and the two r(NN) distances are both longer by 0.040 and 0.169 Å as compared to  $r_e (N_2 {}^1\Sigma_g^+) = 1.0977 \text{ Å}.^{76}$  The fact that the two N-N bond distances differ significantly is in accord with the high covalency of the azido ligand in  $FN_3$ . High covalency facilitates the breakage of the longer  $N_{\alpha}$ - $N_{\beta}$  bond and favors the elimination of  $N_2$ , explaining the high shock sensitivity of these azides. There are three possible binding sites for the addition of  $F^+$  to  $FN_3$ , with the most favorable one being that of  $F^+$  binding to the  $\alpha$ position of FN<sub>3</sub> forming a  $C_s$  structure with a long  $F_2N-N_2$  bond of 2.468 Å and an  $-NF_2$  out-of-plane angle of  $126^{\circ}$  at the CCSD(T)/aVTZ level. The very long N–N bond of 2.468 Å indicates that  $[F_2N-N_2]^+$  would readily lose N2 and, therefore, would be a very unstable species.  $F^+$  addition to the  $\beta$ - and  $\gamma$ -positions form structures of  $C_1$  and  $C_2$  symmetry, which are even less stable by 79.7 and 43.7 kcal/mol, respectively. Similarly,  $F^-$  can bind to  $FN_3$  in three possible positions. Unlike the  $F^+$  addition, the most favorable one is that of  $F^-$  binding to the  $\gamma$ -position of FN<sub>3</sub> forming a planar  $C_s$  structure (C) with N-F bond lengths of 1.608 and 1.508 Å suggesting that this anion might have marginal stability at low temperatures.  $F^-$  addition to the  $\beta$ -position forms structure (B) with  $C_s$  symmetry, and this structure is slightly less stable by 3.0 kcal/mol with a very long N(1)-F(3)bond length of 2.002 Å.  $F^-$  addition to the  $\alpha$ -position leads to a dissociated structure in which F<sup>-</sup> does not bind at the MP2/aVTZ level.

The geometry of FOOF has been reported from microwave spectroscopy<sup>38</sup> and electron diffraction<sup>39</sup> studies, as well as high-level theoretical calculations performed at the CCSD(T)(FC)/CBS<sup>43</sup> and the CCSD(T)/aV5Z<sup>46</sup> levels, both including additional corrections for higher order correlation and core-valence effects. Our MP2/ aVTZ values are in reasonable agreement with the experimental and high-level computational values to within about 0.05 Å with r(O-O) being too short and r(O-F)too long.

The cationic and anionic forms of O<sub>2</sub>F are structurally very different. The  $O_2F^+$  cation is strongly bent ( $\angle OOF =$ 112.9°), has relatively short O-F and O-O bonds, and may have been experimentally observed at low temperatures.<sup>82</sup> The  $O_2F^+$  cation is isoelectronic to the ozone molecule ( $\angle OOO = 117.8^{\circ}$ ),<sup>83</sup> consistent with the predicted bent structure. The first excited  ${}^{3}A''$ , state of  $O_{2}F^{+}$ is only 4.8 kcal/mol higher in energy than the ground  ${}^{1}A$ state at 0 K. In contrast, the  $O_2F^-$  anion is a triplet, almost linear, and is best considered as a very weak complex of  $F^-$  with  $O_2$  with a very long O-F bond distance of 2.844 Å. The triplet state results from binding the closed shell  $F^-$  to the ground state of  $O_2$ , which is a triplet. The first excited singlet state of  $O_2F^-$  is 19.9 kcal/mol higher in energy, consistent with this result and the  ${}^{3}\Sigma_{g}^{-}-{}^{1}\Delta_{g}$  splitting of 22.6 kcal/mol in diatomic O<sub>2</sub>.<sup>76</sup> The neutral radical O<sub>2</sub>F has a C<sub>s</sub> structure<sup>43,84</sup> with a bond angle similar to that of the  $O_2F^+$  cation. The O–O and Q-F bond distances in the neutral O<sub>2</sub>F radical are 0.07 Å and 0.18 Å, respectively, longer than in the cation, suggesting that the radical could be less stable than the cation. The  $r_e$  of  $O_2$  ( ${}^{3}\Sigma_{g}^{-}$ ) is 1.2075 Å<sup>76</sup> and falls in between those of  $O_2F^+$  (1.129 Å) and  $O_2F^-$  (1.228 Å).

The structure of OF<sub>2</sub> is of  $C_{2\nu}$  symmetry, and r(OF) is slightly longer by 0.040 and 0.049 Å than r(OF) of FO  $(^{2}\Pi)$  at the CCSD(T)/aVTZ level<sup>43</sup> and the experimental value,<sup>85</sup> respectively, but in good agreement with the experiment for  $OF_2$ .<sup>86</sup> The addition of  $F^+$  to  $OF_2$  results in a structure of  $C_{3v}$  symmetry with r(OF) slightly shorter than that in OF<sub>2</sub> by 0.027 Å at the CCSD(T)/aVTZ level, suggesting that  $OF_3^+$  might be a marginally stable cation. Addition of  $F^-$  to  $OF_2$  results in a planar T-shaped structure of  $C_{2v}$  symmetry with two long r(OF) distances of 1.743 Å indicating that the  $OF_3^-$  anion could readily lose an F<sup>-</sup> anion and therefore be of low stability.

Addition of  $F^+$  and  $F^-$  to  $O_2F_2$  gives the compounds  $O_2F_3^+$  and  $O_2F_3^-$  with both structures having  $C_s$  symmetry. The  $O_2F_3^+$  cation has relatively short and strong O-F bonds of about 1.38 A but a weak O-O bond of 1.653 Å which would limit its thermal stability and could result in dissociation into OF<sub>2</sub> and OF fragments. In contrast, the corresponding  $O_2F_3^-$  anion has a very short O=O double bond of 1.187 Å, but long and very weak O-F bonds of 1.839 and 2.073 A, suggesting a very unstable species which could easily decompose to  $O_2$  and fluorine.

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Table 2. Components	for Calcu	lated Atomiz	ation Energ	gies in k	cal/mol
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reactions	$CBS^{a}$	$\Delta E_{\text{ZPE}}{}^{b}$	$\Delta E_{\rm CV}{}^c$	$\Delta E_{\rm SR}{}^d$	$\Delta E_{\rm SO}^{\ e}$	$\sum D_0(0 \text{ K})^f$
$NF^+ + e^- \rightarrow N + F$	-204.46	2.23	-0.34	-0.15	-0.39	-207.57
$NF \rightarrow N + F$	77.06	1.63	0.06	-0.25	-0.39	74.85
$NF^- \rightarrow N + F + e^-$	81.86	1.07	0.10	-0.26	-0.39	80.24
$OF^+ + e^- \rightarrow O + F$	-241.30	1.85	-0.30	-0.06	-0.61	-244.12
$OF^- \rightarrow O + F + e^-$	104.62	1.08	0.08	-0.23	-0.61	102.78
$F_2^+ + e^- \rightarrow 2F$	-324.57	1.60	-0.41	0.07	-0.78	-327.29
$F_2^- \rightarrow 2F + e^-$	107.35	0.64	0.14	-0.13	-0.78	105.93
$NF_2^+ + e^- \rightarrow N + 2F$	-121.90	4.65	-0.33	-0.41	-0.78	-128.07
$NF_2 \rightarrow N + 2F$	145.07	3.74	0.03	-0.51	-0.78	140.08
$NF_2^- \rightarrow N + 2F + e^-$	169.73	2.76	0.07	-0.51	-0.78	165.76
$HNF_2 \rightarrow H + N + 2F$	226.69	12.26	0.19	-0.61	-0.78	213.23
$NF_3^+ + e^- \rightarrow N + 3F$	-84.80	7.47	0.05	-0.85	-1.17	-94.24
$NF_3 \rightarrow N + 3F$	205.61	6.55	0.01	-0.67	-1.17	197.23
$NF_3^- \rightarrow N + 3F + e^-$	232.39	4.20	0.14	-0.71	-1.17	226.45
$NF_4^+ + e^- \rightarrow N + 4F$	-16.67	10.19	-0.02	-1.10	-1.56	-29.55
$NF_4 \rightarrow N + 4F$	204.33	6.70	0.02	-0.89	-1.56	195.20
$NF_4^- \rightarrow N + 4F + e^-$	298.82	6.01	-0.19	-0.84	-1.56	281.22
$NH_5 \rightarrow N + 5H$	300.51	27.38	0.52	-0.21	0.00	273.43
$NF_5 \rightarrow N + 5F$	205.76	9.94	-0.10	-1.32	-1.95	192.45
$N_2F^+ + e^- \rightarrow 2N + F$	-42.70	5.46	0.61	-0.36	-0.39	-47.70
$N_2F \rightarrow 2N + F$	227.79	3.49	0.65	-0.16	-0.39	224.40
$N_2F^- \rightarrow 2N + F + e^-$	306.50	3.55	0.75	-0.34	-0.39	306.50
$cis-N_2F_2 \rightarrow 2N + 2F$	250.89	7.27	0.38	-0.69	-0.78	242.53
<i>trans</i> -N <sub>2</sub> F <sub>2</sub> $\rightarrow$ 2N + 2F	249.45	7.12	0.40	-0.71	-0.78	241.25
$TS-N_2F_2 \rightarrow 2N + 2F$	180.82	5.78	0.71	-0.75	-0.78	174.22
$N_2F_2(^{3}B) \rightarrow 2N + 2F$	189.02	4.97	0.27	-0.69	-0.78	182.85
$F_2N = N \rightarrow 2N + 2F$	236.31	6.97	0.68	-0.73	-0.78	228.51
$N_2F_3^+ + e^- \rightarrow 2N + 3F$	24.98	10.54	0.53	-0.90	-1.17	12.89
$N_2F_3 \rightarrow 2N + 3F$	251.70	8.82	0.31	-0.88	-1.17	241.13
$N_2F_3 \rightarrow 2N + 3F + e^-$	339.28	7.60	0.49	-0.91	-1.17	330.09
$N_2F_4(C_{2h}) \rightarrow 2N + 4F$	314.36	11.25	0.19	-1.11	-1.56	300.63
$N_2F_4(C_2) \rightarrow 2N + 4F$	314.16	11.27	0.26	-1.15	-1.56	300.45
$N_2F_5^+ + e^- \rightarrow 2N + 5F$	104.90	12.51	0.09	-1.25	-1.95	89.28
$N_2F_5^- \rightarrow 2N + 5F + e^-$	380.48	9.89	0.02	-1.35	-1.95	367.31
$FN_3 \rightarrow 3N + F$	283.76	7.91	1.00	-0.67	-0.39	275.79
$F_2N_3^+(A, C_s) + e^- \rightarrow 3N + 2F$	112.15	9.03	0.30	-0.62	-0.78	102.02
$F_2N_3^+$ ( <b>B</b> , $C_1$ ) + $e^- \rightarrow 3N + 2F$	33.68	9.90	0.21	-0.87	-0.78	22.34
$F_2N_3^+(C, C_2) + e^- \rightarrow 3N + 2F$	68.95	10.07	0.67	-0.43	-0.78	58.33
$F_2N_3^-$ ( <b>B</b> , $C_s$ ) $\rightarrow$ 3N + 2F + e <sup>-</sup>	383.94	9.21	0.99	-0.86	-0.78	378.01
$F_2N_3^-(C, C_s) \rightarrow 3N + 2F + e^-$	386.44	8.51	0.84	-0.94	-0.78	377.05
$OF_2 + e^- \rightarrow O + 2F$	92.66	3.24	-0.10	-0.31	-1.00	88.02
$OF_3^+ \rightarrow O + 3F + e^-$	-193.54	5.42	-0.54	-0.19	-1.39	-201.08
$OF_3^- \rightarrow O + 3F + e^-$	188.37	3.65	-0.23	-0.42	-1.39	182.68
$O_2F^+({}^{1}A', C_s) + e^- \rightarrow 2O + F$	-152.13	3.90	-0.31	-0.11	-0.83	-157.28
$O_2F^+$ ( $^{s}A'', C_s$ ) + $e^- \rightarrow 2O + F$	-155.25	5.64	-0.10	-0.22	-0.83	-162.04
$O_2F^-(A'', C_s) \rightarrow 2O + F + e^-$	201.19	2.18	0.30	-0.46	-0.83	198.02
$O_2F^-(^{1}A', C_s) \rightarrow 2O + F + e^-$	182.46	2.94	0.09	-0.59	-0.83	178.19
$O_2F_2 \rightarrow 2O + 2F$	149.90	5.08	-0.10	-0.46	-1.22	143.04
$O_2F_3^+ + e^- \rightarrow 2O + 3F$	-140.69	6.77	-0.60	-0.34	-1.61	-150.01
$O_2F_3 \rightarrow 2O + 3F + e^-$	256.53	8.08	-0.12	-0.65	-1.61	246.07

<sup>*a*</sup> Extrapolated by using eq 1 with the aVnZ, n = D, T, Q basis set. <sup>*b*</sup> The zero point energies were obtained as described in the text. <sup>*c*</sup> Core-valence corrections were obtained with the cc-pwCVTZ (N, O, and F) at the optimized CCSD(T) or MP2 geometries. <sup>*d*</sup> The scalar relativistic correction is based on a CISD(FC)/VTZ MVD calculation and is expressed relative to the CISD result without the MVD correction. <sup>*e*</sup> Correction due to the incorrect treatment of the atomic asymptotes as an average of spin multiplets. Values are based on C. Moore's Tables, ref 71. <sup>*f*</sup> The theoretical value of the dissociation energy to atoms  $\sum D_0(0 \text{ K})$ .

**Frequencies.** The anharmonic vibrational frequencies of NF<sub>2</sub>,<sup>11,12</sup> HNF<sub>2</sub>,<sup>13</sup> NF<sub>3</sub>,<sup>13</sup> *cis*-N<sub>2</sub>F<sub>2</sub>,<sup>16</sup> and *trans*-N<sub>2</sub>F<sub>2</sub>,<sup>13,15</sup> N<sub>2</sub>F<sub>4</sub>,<sup>19</sup> NF<sub>4</sub><sup>+</sup>,<sup>14</sup> N<sub>2</sub>F<sub>3</sub><sup>+</sup>,<sup>18</sup> and N<sub>2</sub>F<sup>+17</sup> have been experimentally observed, and our current and previously calculated CCSD(T)/aVTZ harmonic frequencies are in excellent agreement with them (see Supporting Information). The largest discrepancies are found for the F–N=N bend ( $\pi$  mode) of N<sub>2</sub>F<sup>+</sup> which was calculated to be 20 cm<sup>-1</sup> higher than the experimental value and the N–H stretch of HNF<sub>2</sub> which was calculated to be 175 cm<sup>-1</sup> (a' mode) larger than the experimental anharmonic value of 3193 cm<sup>-1</sup>; the latter difference is due to the large anharmonic vibrational frequencies of FN<sub>3</sub>

have also been reported,<sup>20</sup> and our calculated CCSD(T)/ VTZ harmonic frequencies are in very good agreement, with a maximum discrepancy of 33 cm<sup>-1</sup>, except for the N=N stretching vibration (a' mode) which is predicted to be larger by 62 cm<sup>-1</sup> than the reported gas-phase experimental value of 2037 cm<sup>-1</sup>. We note that there is a small imaginary frequency at the CCSD(T)/aT level for NF<sub>4</sub><sup>-</sup>, but that at the analytic MP2/aT level, the frequencies are all real for the  $C_{2\nu}$  structure.

Heats of Formation. The energetic components for predicting the total molecular dissociation energies are given in Table 2. We first discuss some general trends in the atomization energy components. The  $\Delta E_{\rm CV}$  corrections are small and may be positive or negative with most

being less than |1.0| kcal/mol. The  $\Delta E_{\rm SR}$  corrections are all small and negative, ranging from -0.06 (OF<sup>+</sup>) to -1.35 (N<sub>2</sub>F<sub>5</sub><sup>-</sup>) kcal/mol, except for F<sub>2</sub><sup>+</sup> (0.07 kcal/mol). We estimate that the error bars for the calculated heats of formation are  $\pm 1.5$  kcal/mol (except for FOOF), considering errors in the energy extrapolation, frequencies, and other electronic energy components. An estimate of the potential for significant multireference character in the wave function can be obtained from the  $T_1$  diagnostic<sup>87</sup> for the CCSD calculation. The values for the  $T_1$  diagnostics (Supporting Information, Table SI-4) are small showing that the wave functions are dominated by single configurations.

The calculated heats of formation at 0 and 298 K are given in Table 3, and we use the values at 298 K in our discussion below. The heats of formation of NF<sub>2</sub> and NF<sub>3</sub> have been reported in the NIST-JANAF<sup>8</sup> tables, and our calculated values are in excellent agreement within 1.3 kcal/mol and within the reported error bars. Our calculated value for the heat of formation of NF is predicted to be 3.4 kcal/mol more stable than the JANAF value of 59.5 kcal/mol and within the  $\pm 7.9$  kcal/mol error bars.<sup>8</sup> Our calculated value for the heat of formation of  $N_2F_4$ is 2.0 kcal/mol more stable than the JANAF value of -2.0 kcal/mol,<sup>8</sup> and within the  $\pm 2.5$  kcal/mol error bars. Our composite CCSD(T)/CBS value for the  $\Delta H_{f,298K}$ -(FOOF) is given simply because we needed it for the F and F<sup>-</sup> affinities. It differs from the experimental value adopted by the NIST-JANAF tables of  $4.59 \pm 0.5$  kcal/mol<sup>8</sup>

by 5 to 6 kcal/mol and from the best calculated value of  $6.4 \pm 0.7$  kcal/mol by 4.5 kcal/mol, and the reasons for this have been discussed in detail.<sup>43</sup> The structure, bonding, harmonic vibrational frequen-

cies, and decomposition reactions of NF5 have been predicted by Bettinger et al.<sup>81</sup> employing various correlated levels of theory up through the CCSD(T) level with basis sets of triple- $\zeta$  quality. They predict the overall reaction of  $NF_5 \rightarrow NF_3 + F_2$  to be exothermic by 42.2 kcal/mol at the CCSD(T)/VTZ//CCSD/DZP level, in excellent agreement with our higher level value of 41.7 kcal/mol. We predict the decomposition of NF<sub>5</sub> into NF<sub>4</sub>  $(C_{3v}$  symmetry) and an F radical to be exothermic by 2.7 kcal/mol at the CCSD(T)/CBS//MP2/aVTZ level plus additional corrections (Table 5). The value of 8.5 kcal/ mol obtained at the CCSD(T)/VTZ//CCSD/DZP + ZPVE level<sup>81</sup> is in qualitative agreement with our higher level value. The NF<sub>4</sub> radical is predicted to be unbound with respect to  $NF_3 + F$  by 2.0 kcal/mol (Table 5) at the CCSD(T)/CBS level plus additional corrections, whereas Bettinger et al.<sup>81</sup> find it to be weakly bound by only 0.1 kcal/mol at the CCSD(T)/VTZ//CCSD/DZP + ZPVE level. Thus, NF<sub>5</sub> would also readily decompose to NF<sub>3</sub> with the generation of two fluorine radicals in an exothermic reaction (-4.8 kcal/mol). Unlike the analogous fluorine counterpart, the trigonal bipyramidal form of  $NH_5$  ( $D_{3h}$ ) is not predicted to be a minimum at the CCSD(T)/aVTZ level with an imaginary frequency of 982.4 cm<sup>-1</sup> ( $a_2''$  mode). Given  $\Delta H_{f,0K}(NH_3)^{8,88}$  the decomposition pathway of  $NH_5 \rightarrow NH_3 + H_2$  is predicted to be highly exothermic by 106.9 kcal/mol which is greater

Table 3. Calculated Heats of Formation (kcal/mol) at 0 and 298 K

molecule	$\Delta H_f(0 \text{ K})_{\text{theory}}$	$\Delta H_f$ (298 K) <sub>theory</sub>	$\Delta H_f(298 \text{ K})_{exp}$
$\mathrm{NF}^+(^2\Pi, C_{\mathrm{ouv}})$	338.6	338.6	
NF $(^{3}\Sigma^{-}, C_{\infty v})$	56.1	56.1	$59.5 \pm 7.9^{8}$
$NF^{-}(^{2}\Pi, C_{\infty v})$	50.8	50.8	
$OF^+(^3\Sigma^-, C_{\infty\nu})$	321.6	321.6	
$\mathrm{OF}^{-}(^{1}\Sigma^{+}, C_{\infty v})$	-25.3	-25.3	
$F_2^+(^2\Pi_g, D_{\infty h})$	364.2	364.2	
$F_2^{-}({}^1\Sigma^+_{g}, D_{\infty h})$	-69.0	-69.0	
$NF_2^+({}^1A_1, C_{2\nu})$	277.5	276.8	
$NF_2(^2B_1, C_{2\nu})$	9.4	8.8	$10.1 \pm 1.9^{8}$
$NF_2^{-}(^1A_1, C_{2\nu})$	-16.3	-16.8	
$HNF_2(^1A', C_s)$	-12.1	-13.7	
$NF_3^+(^2A_1, C_{3\nu})$	262.2	260.8	
$NF_3(^{1}A_1, C_{3v})$	-29.3	-30.7	$-31.6 \pm 0.3^{8}$
$\mathrm{NF_3}^-(^2\mathrm{A}',C_s)$	-58.5	-58.9	
$\mathrm{NF_4}^+(^1\mathrm{A_g}, T_d)$	216.0	213.8	
$NF_4(^2A_1, C_{3v})$	-8.8	-9.6	
$NF_4^{-}({}^{1}A_1, C_{2\nu})$	-94.5	-93.4	
$\rm NH_5({}^1A_1{}', D_{3h})$	97.3	93.7	
$NF_5({}^1A_1', D_{3h})$	12.4	10.1	
$N_2F^+(^1\Sigma_g, C_{\infty v})$	291.2	290.6	
$N_2 F(^2\Sigma_g, C_{\infty v})$	19.1	19.5	
$N_2F^-({}^1A', C_s)$	-63.0	-62.8	
$cis-N_2F_2(^1A_1, C_{2v})$	19.5	18.1	$16.4 \pm 1.2^{8}$
trans-N <sub>2</sub> F <sub>2</sub> ( $^{1}A_{g}, C_{2h}$ )	20.8	19.5	$19.4 \pm 1.2^{8}$
$\mathrm{TS-N_2F_2}(^1\mathrm{A}',C_s)$	87.8	86.8	
$N_2F_2(^{3}B, C_2)$	79.1	78.0	
$\mathbf{F}_2 N = \mathbf{N} \left( {}^1 \mathbf{A}_1, C_{2\nu} \right)$	33.5	32.4	
$N_2F_3^+({}^1A', C_s)$	267.6	265.6	
$N_2F_3$ ( <sup>2</sup> A, C <sub>1</sub> )	39.3	37.6	
$N_2F_3^{-}(^1A, C_1)$	-49.6	-50.8	
$N_2F_4$ ( <sup>1</sup> A <sub>g</sub> , C <sub>2h</sub> )	-1.7	-4.0	$-2.0 \pm 2.5^{8}$
$N_2F_4$ ( <sup>1</sup> A, C <sub>2</sub> )	-1.5	-3.9	
$N_2F_5^+({}^1A', C_s)$	228.1	230.6	
$N_2F_5^{-}(^1A, C_1)$	-49.9	-51.1	
$FN_3(^1A', C_s)$	80.3	79.0	
$\mathbf{F}_{2}\mathbf{N}_{3}^{+}(^{1}\mathbf{A}^{\prime},C_{s},\mathbf{A})$	272.5	271.8	
$F_2N_3^+$ ( <sup>1</sup> A, C <sub>1</sub> , <b>B</b> )	352.2	350.4	
$F_2N_3^+$ ( <sup>1</sup> A, C <sub>2</sub> , C)	316.2	314.5	
$F_2N_3^{-}(^1A', C_s, \mathbf{B})$	0.5	-1.1	
$F_2N_3^{-}(^1A', C_s, \mathbf{C})$	-2.5	-4.0	
$OF_2(^1A_1, C_{2\nu})$	7.9	7.4	
$OF_3^+(^{1}A_1, C_{3\nu})$	315.5	314.2	
$OF_3^{-}(^{1}A_1, C_{2\nu})$	-68.3	-68.9	
$O_2F^+$ ( <sup>1</sup> A', $C_s$ )	293.7	293.3	
$O_2F^+({}^{3}A'', C_s)$	298.5	298.4	
$O_2 F^- ({}^3 A^{\prime\prime}, C_s)$	-61.6	-62.3	
$O_2 F^-(^1 A', C_s)$	-41.7	-42.0	0
$O_2F_2(^{1}A, C_2)$	11.9	10.9	$4.6 \pm 0.5^{8}$
			$5.9 \pm 0.4^{9}$
			$4.6 \pm 0.2^{9}$
			$4.7 \pm 0.3^{10}$
$O_2F_3^+({}^{'}A', C_s)$	323.4	321.4	
$O_2F_3^{-}({}^{'}A', C_s)$	-72.7	-74.0	

than the bond energy<sup>8</sup> in  $H_2$ , so two H atoms will be readily formed.

**Ionization Potentials and Electron Affinities.** The ionization potentials (IP) and electron affinities (EA) for a number of the molecules are given in Table 4. For the diatomics, the agreement with experimental data<sup>89–93</sup> is

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**Table 4.** Calculated Ionization Potentials (IP) and Electron Affinities (EA) of the  $N_x F_v$  and  $O_x F_v$  Compounds in eV at 0 K

molecule	IP (eV)	IP expt	EA eV	EA expt
F <sub>2</sub>	15.79	$15.697 \pm 0.003^{90}$	2.99	$3.005 \pm 0.071^{91}$
FO	12.80	$12.77^{88}$	2.24	$2.2720 \pm 0.0060^{89}$
NF	12.25	$12.26 \pm 0.01^{92}$	0.23	$0.50^{57}$
$NF_2$	11.63	$11.63 \pm 0.01^{93}$	1.11	$1.21 \pm 0.20^{94}$
				$1.10 \pm 0.10^{95}$
NF <sub>3</sub>	12.64	$12.94 \pm 0.01^{97}$	1.27	
NF <sub>4</sub>	9.75		3.72	
$N_2F$	11.80		3.56	
$N_2F_3$	9.90		3.86	
trans-N <sub>2</sub> F <sub>2</sub>	12.69	12.896		
$cis-N_2F_2$	13.28			

good with the largest error in the ionization potential being 0.1 eV for  $F_2$  and the largest error in the electron affinity being  $\sim 0.25$  eV for NF. Our calculated value for EA(NF) is in excellent agreement within 0.02 eV of the value of 0.247 eV reported using even larger basis sets with double diffuse functions.<sup>56</sup> There are no experimental details<sup>57</sup> reported for the experimental value of 0.5 eV so this value needs to be remeasured. The difference in the IP of F<sub>2</sub> is most likely due to the difference in the vertical and adiabatic values. At the CCSD(T)/CBS level with the aug-cc-pVnZ basis sets for n = D, T, Q, the difference in the vertical and adiabatic IPs ( $\Delta E_{elec}$ ) is 0.17 eV, consistent with this argument. This suggests that there may be differences in the calculated and experimental IPs because the calculated value is an adiabatic value and the experimental value is an extrapolation of a vertical value to the adiabatic value.

The calculated adiabatic IP and EA for  $NF_2$  are in excellent agreement with experiment.<sup>94–96</sup> The calculations show that the predicted adiabatic IP for trans-FNNF agrees within 0.1 eV with the experimental value<sup>97</sup> of 12.8 eV which is only given to one decimal place accuracy. The calculated IP of 13.28 eV of the  $cis-N_2F_2$ is substantially higher. This difference in the ionization potentials for the cis and trans isomers is considerably larger than those in *cis*- and *trans*-1,2-difluoroethylene where the IPs are within 0.02 eV of each other.<sup>98</sup> The calculated IP for NF<sub>3</sub> is 0.3 eV below the experimental value.99 Examination of the spectra shows a broad photoionization peak. We calculated the  $\Delta E_{elec}$  component of the difference for the vertical and adiabatic ionization potentials of NF<sub>3</sub> at the CCSD(T)/CBS level and predict a substantial difference of 1.19 eV. The large difference in the adiabatic and vertical IPs is due to the substantial changes in the geometry of  $\sim 0.1$  A in r(N-F) and of  $\sim 12^{\circ}$ in the F-N-F bond angle. We further checked this by calculating the ionization potential of NH<sub>3</sub>. The difference in the adiabatic and vertical values is 0.74 eV. The calculations show that it is very difficult to extrapolate from the vertical to the adiabatic value for  $NF_3$ . At the CCSD(T)/CBS level, the vertical and adiabatic ionization potentials for cis- and trans-N<sub>2</sub>F<sub>2</sub> are also predicted to have large differences of 0.77 and 0.82 eV, respectively.

The electron affinity of NF<sub>3</sub> of 1.27 eV is surprisingly large for a closed shell molecule. The electron affinities of NF<sub>4</sub>, N<sub>2</sub>F, and N<sub>2</sub>F<sub>3</sub> are all comparable and up to 0.5 eV larger than that of 3.40 eV for the F atom.<sup>8</sup>

**A**–**F** Bond Dissociation Energies. The adiabatic bond dissociation energies (BDEs) at 0 K, in which dissociation

Table 5. Calculated A–F Bond Dissociation Energies (BDE) of the  $N_x F_y$  and  $O_x F_y$  Compounds in kcal/mol at 0 K

molecule	product	BDE
NF	N + F	74.9
NF <sub>2</sub>	NF + F	65.2
NF <sub>3</sub>	$NF_2 + F$	57.2
NF <sub>4</sub>	$NF_3 + F$	-2.0
NF <sub>5</sub>	$NF_4 + F$	-2.7
N <sub>2</sub> F	$N_2 + F$	-0.6
$N_2F_2$	$N_2F + F$	18.1
$F_2N=N$	$\overline{N_2F} + F$	4.1
$N_2F_3$	$\overline{N_2F_2} + F$	57.8
$N_2F_4$	$\overline{N_2F_3} + F$	59.5
OF	O + F	51.2
$OF_2$	OF + F	36.9

of the reactant occurs to the ground states of the separated product species, are given in Table 5. The series NF,  $NF_2$ ,  $NF_3$  exhibits a moderate substituent effect. The addition of one F atom results in a decrease of the BDE from 74.9 kcal/mol in NF to 57.2 kcal/mol in NF<sub>3</sub>. NF<sub>4</sub> and NF<sub>5</sub> are predicted to have negative bond energies and are metastable species as discussed above. For N<sub>2</sub>F, there are two possible dissociation pathways with the formation of  $N_2 + F$  being favored by 150 kcal/mol over that resulting in NF + N. The N–N BDE in N<sub>2</sub>F is 149.6 kcal/mol, about 75 kcal/mol less than that in  $N_2$ .<sup>44,76</sup> The N-F BDE in  $N_2F_2$  is predicted to be 18.1 kcal/mol, that is, to be 57 kcal/mol lower than that in diatomic N-F. The N-F BDE in the  $F_2N=N$  isomer has a small value of 4.1 kcal/mol and is 14.0 kcal/mol less than that in  $N_2F_2$ . The N-F BDEs in  $N_2F_3$  and  $N_2F_4$  are predicted to have moderately high values of about 59 kcal/mol, to be within 2 kcal/mol of each other, are bracketed by the N-F BDEs of NF<sub>2</sub> and NF<sub>3</sub>, and are 17.1 and 15.4 kcal/mol less, respectively, than that of diatomic NF. The O-F BDEs in OF<sub>2</sub> and OF have intermediate values of 36.9 and 51.2 kcal/mol, respectively, and the BDE of OF<sub>2</sub> is predicted to be 14.3 kcal/mol lower than that in diatomic OF.<sup>45</sup>

 $\mathbf{F}^+$  and  $\mathbf{F}^-$  Affinities. Given the experimental  $\Delta H_{f,0K}$ -( $\mathbf{F}^+$ ) = 419.40 and  $\Delta H_{f,0K}(\mathbf{F}^-)$  = -59.96 kcal/mol, we can predict the  $\mathbf{F}^+$  and  $\mathbf{F}^-$  affinities of the N<sub>x</sub>F<sub>y</sub> and O<sub>x</sub>F<sub>y</sub> compounds defined as  $-\Delta H$  for the representative reactions of N<sub>2</sub>F<sub>y</sub>:

$$N_2F_y + F^+ \rightarrow N_2F_{y+1}^+ \tag{3}$$

$$N_2F_y + F^- \rightarrow N_2F_{y+1}^- \tag{4}$$

The calculated  $F^+$  and  $F^-$  affinities at 0 K are given in Table 6. The  $F^+$  cation affinity (FCA) within the N<sub>2</sub>F<sub>y</sub> series increases from N<sub>2</sub> to N<sub>2</sub>F<sub>4</sub> with the FCA of N<sub>2</sub>F<sub>2</sub> being 44 kcal/mol higher than that of N<sub>2</sub>, and FCA(N<sub>2</sub>F<sub>4</sub>) being 18 kcal/mol larger than FCA(N<sub>2</sub>F<sub>2</sub>). The FCA-(N<sub>2</sub>F<sub>2</sub>) and FCA(N<sub>2</sub>F<sub>4</sub>) are toward the more positive end of the previously reported oxidizer strength scale<sup>1</sup> with values comparable to those of XeF<sub>2</sub>O (173 kcal/mol) and ClFO (193 kcal/mol), respectively, indicating that the corresponding cations are only moderately strong oxidizers. The FCA(N<sub>2</sub>F<sub>4</sub>) is 4 kcal/mol lower and the FCA-(NF) is 4.6 kcal/mol higher than that of the N atom. The FCA(N<sub>2</sub>) falls in the less positive range of the oxidizer scale<sup>1</sup> with a value that is bracketed by FCA(O<sub>2</sub>)

Table 6. Calculated  $F^+$  and  $F^-$  Affinities of the  $N_xF_y$  and  $O_xF_y$  Compounds in kcal/mol at 0 K

molecule	$\mathrm{F}^+$	$F^{-}$
N	193.4	1.8
0	156.8	24.3
F	73.6	27.5
NF	198.0	12.5
NF <sub>2</sub>	166.6	7.9
NF <sub>3</sub>	174.1	5.2
N <sub>2</sub>	127.0	-1.5
cis-N <sub>2</sub> F <sub>2</sub>	171.3	9.1
$trans-N_2F_2$	172.6	10.4
FN <sub>3</sub>	227.2	22.8
$N_2F_4$	189.6	-11.8
02	125.7	1.6
OF <sub>2</sub>	111.8	16.2
$O_2 \tilde{F_2}$	107.9	24.6

(126 kcal/mol) and FCA(BrF<sub>3</sub>O) (131 kcal/mol). Therefore, N<sub>2</sub>F<sup>+</sup> is a relatively strong oxidative fluorinator. The FCA(NF<sub>2</sub>) and FCA(NF<sub>3</sub>) are comparable to FCA-(N<sub>2</sub>F<sub>2</sub>) and are in the more positive end of the oxidizer strength scale,<sup>1</sup> making NF<sub>3</sub><sup>+</sup> and NF<sub>4</sub><sup>+</sup> only moderately strong oxidizers. With FCA(FN<sub>3</sub>) of 227.2 kcal/mol,  $F_2N_3^+$  is predicted to be the weakest oxidizer of the compounds studied.

Although FCA(O<sub>2</sub>) and FCA(N<sub>2</sub>) are very similar and differ only by 1.3 kcal/mol, FCA(O<sub>2</sub>F<sub>2</sub>) is very different from FCA(N<sub>2</sub>F<sub>2</sub>), with O<sub>2</sub>F<sub>3</sub><sup>+</sup> being a considerably stronger oxidizer than N<sub>2</sub>F<sub>3</sub><sup>+</sup>. The oxidizing power of O<sub>2</sub>F<sub>3</sub><sup>+</sup> also exceeds that of O<sub>2</sub>F<sup>+</sup> by 18 kcal/mol, contrary to the large decrease of 44 kcal/mol in oxidizing power predicted for going from N<sub>2</sub>F<sup>+</sup> to N<sub>2</sub>F<sub>3</sub><sup>+</sup>. The oxidizing power of OF<sub>3</sub><sup>+</sup> falls between those of O<sub>2</sub>F<sup>+</sup> and O<sub>2</sub>F<sub>3</sub><sup>+</sup>, while OF<sup>+</sup> is predicted to be the weakest oxidizer among the group of oxygen fluoride cations studied.

Neither  $N_2$  nor  $N_2F_4$  is predicted to bind F<sup>-</sup>, and for  $N_2$ , this is consistent with the expected interaction between the closed-shell atomic anion F<sup>-</sup> and the very stable  $N_2$  molecule. As would be expected as well, there is no place for F<sup>-</sup> to interact with N<sub>2</sub>F<sub>4</sub> without encountering an electron pair, so it too does not bind  $F^-$ . N<sub>2</sub>F<sub>2</sub> is a very weak Lewis acid with an  $F^-$  affinity (FA) of about 10 kcal/mol for either the cis or trans isomers; structurally,  $N_2F_3^-$  is a weakly bound complex of  $N_2F_2$  and  $F^-$ .  $NF_2$ and NF3 are also predicted to be very weak Lewis acids, as is NF with FA(NF)  $\sim$  13 kcal/mol. The FA(FN<sub>3</sub>) is the largest compared to the other  $N_x F_v$  compounds studied and at 23 kcal/mol is almost two times that of NF. However, its value is still too low for the probable existence of a room-temperature stable  $N_3F_2^-$  salt. The  $FA(NF_3)$  of 5.2 kcal/mol is one of the smallest for the NF<sub>x</sub> compounds, and 2.7 kcal/mol smaller than  $FA(NF_2)$  and only 3.4 kcal/mol larger than FA(N).

Unlike N<sub>2</sub>, O<sub>2</sub> is predicted to have a small, but positive  $F^-$  affinity, indicating a weak complex of O<sub>2</sub> with  $F^-$  that is slightly more stable than the separated O<sub>2</sub> +  $F^-$  reactants. The  $F^-$  affinities of O<sub>2</sub> and N are essentially the same, but FA(O) is predicted to be 23 kcal/mol higher than those of O<sub>2</sub> and N. The FA(O<sub>2</sub>F<sub>2</sub>) is positive and its

Lewis acidity increases by 23 kcal/mol upon  $F_2$  addition to  $O_2$ .  $OF_2$  is a very weak Lewis acid and its FA falls in between  $FA(O_2)$  and  $FA(O_2F_2)$ . Of all the molecules and atoms of this study, F has the highest  $F^-$  affinity, slightly above those of O and  $O_2F_2$ , but none can be considered to exhibit significant Lewis acidity.

NF<sub>4</sub> Radical. The lowest unoccupied molecular orbital (LUMO) of NF<sub>4</sub><sup>+</sup> ( $T_d$  symmetry) is an  $a_1$  orbital so the addition of an electron to give neutral  $NF_4$  does not necessarily lead to a distortion. The equivalent  $C_{2\nu}$  structure of NF<sub>4</sub> is not a minimum based on a harmonic vibrational analysis at the MP2/aVTZ level (one imaginary frequency of 628.1 cm<sup>-1</sup> (b<sub>1</sub> mode)). A search for the minimum of NF<sub>4</sub> led to a  $C_{3\nu}$  structure with a long N-F bond of 2.702 Å along the  $C_3$  axis and an N-F distance in the NF<sub>3</sub> moiety of 1.369 Å, 0.002 Å shorter than the N-F distance in NF<sub>3</sub> (<sup>1</sup>A<sub>1</sub>) at the MP2/aVTZ level. The  $C_{3\nu}$ structure of NF4 is best described as a loose complex of an F atom weakly interacting with NF<sub>3</sub>, and is predicted to be minimum at the MP2/aVTZ level with two low frequencies of 15.3 cm<sup>-1</sup> (e mode), corresponding to an N-F wag, and 56.0 cm<sup>-1</sup> ( $a_1$  mode), corresponding to the unique N-F stretch. The IP of NF<sub>4</sub> is predicted to be 9.75 eV, substantially lower than that in other NF compounds consistent with the formation of a more stable  $NF_4^+$  ion. Bronsted Acidity of HNF<sub>2</sub> and the Elusive NF<sub>2</sub><sup>-</sup> Anion.

The gas phase acidity  $(HNF_2 \rightarrow NF_2^- + H^+)$  can be predicted  $(\Delta H_f(H^+, 0 \text{ K}) = 365.22 \text{ kcal/mol})^{8,100}$  and the enthalpic contribution to the deprotonation energy is 361.1 kcal/mol. There have been many unsuccessful attempts to isolate the corresponding  $NF_2^-$  anion.<sup>101</sup> A likely explanation for these failures is the ease with which  $NF_2^-$  can lose an  $F^-$  anion to form NF, which can then readily dimerize in a highly exothermic reaction (92.8 kcal/mol at the CCSD(T)/CBS level plus additional corrections) to form  $N_2F_2$  with the formation of an N=N double bond. We predict the enthalpy of the reaction  $NF_2^- \rightarrow NF + \hat{F}^-$  (-FA(NF), Table 6) to be only 12.5 kcal/mol at 0 K, showing that NF<sub>2</sub><sup>-</sup> is only marginally stable toward  $F^-$  loss. The possibility of  $NF_2^$ dissociating to NF<sup>-</sup> + F radicals is much less likely because this pathway is highly endothermic by 85.5 kcal/mol.

#### Conclusions

We have predicted the heats of formation of a number of small neutral and ionic  $N_xF_y$  and  $O_xF_y$  systems at the CCSD(T) level plus additional corrections. The calculated heats of formation and stabilities are in good agreement with the available experimental values, except for FOOF because of the exclusion of higher order correlation effects in this

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Figure 2. Summary of structures of the stable cations and unstable anions.

molecule. The cations are in general more stable than the anions (Figure 2). The  $F^+$  affinities of the  $N_2F_y$  molecules increase by 44 kcal/mol from  $N_2$  to  $N_2F_4$  whereas those of the  $O_2F_y$  molecules exhibit an opposite trend with a notable decrease of 18 kcal/mol from  $O_2$  to  $O_2F_2$ . Neither  $N_2$  nor  $N_2F_4$  are predicted to bind  $F^-$ , whereas  $N_2F_2$  is a very weak Lewis acid with an  $F^-$  affinity of about 10 kcal/mol for either the cis or trans isomer.  $O_2$  is predicted to have a small but

positive  $F^-$  affinity while that of  $O_2F_2$  is also positive and increases by 23 kcal/mol upon  $F_2$  addition to  $O_2$ .

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**Supporting Information Available:** Total CCSD(T) energies (Table SI-1) as a function of basis set; calculated CCSD(T)/aVTZ and MP2/aVTZ frequencies (cm<sup>-1</sup>) (Table SI-2); components for calculated reactions energies in kcal/mol (Table SI-3); and  $T_1$  diagnostics calculated at the CCSD(T)/aVQZ level (Table SI-4). This material is available free of charge via the Internet at http://pubs.acs.org.